

On the l -function and the Reduced Volume of Perelman II *

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1 Introduction

This paper is a sequel of [Y2]. In [Y2], a number of geometric and analytic properties of the l -function and the reduced volume of Perelman were derived. In this sequel, we present a major application of the l -function and the reduced volume, namely their application to the analysis of the asymptotical limits of κ -solutions. Our focus is to present a complete and detailed proof of Proposition 11.2 in [P1], which provides asymptotical convergence of κ -solutions and identifies the asymptotical limits to be nonflat gradient shrinking solitons. A main foundation for the proof is a study of the limit l -function and the asymptotical reduced volume (see Definition 1) on the asymptotical limits of κ -solutions. This is an extension and further development of the theory presented in [Y2].

Recall that κ -solutions are κ -noncollapsed nonflat ancient solutions of the Ricci flow with bounded and nonnegative curvature operator. In dimension three, rescaling limits of a solution of the Ricci flow near a blow-up singularity are κ -solutions. In higher dimensions they also arise as the same kind of rescaling limits under suitable assumptions on the curvature. In other words, in various situations, blow-up singularities of the Ricci flow are modelled on κ -solutions. For this reason, it is important to understand the structures of κ -solutions. To this purpose, one performs blow-down rescalings of a given κ -solution, extracts smooth limits via the process, and identifies them. This is precisely the content of [Proposition 11.2, P1]. This proposition plays an important role in the analysis of blow-up singularities given in Perelman's papers [P1] and [P2].

Now we give a more detailed account of the content of this paper. In Section 2, we first derive an upper bound for Δl (the Laplacian of the l -function) in the weak sense. Then we derive an analytic lemma concerning strong convergence of Sobolev

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functions. These two tools are important for proving [Proposition 11.2, P1], as will be explained below.

Section 3 and Section 4 are devoted to the proof of [Proposition 11.2, P1]. The proof follows the sketch of proof in [P1] for [Proposition 11.2, P1] and incorporates the ideas in Sections 9 and 10 in [P1] regarding how to identify gradient shrinking solitons. The statements of the theorem and the proof of the convergence part are given in Section 3, while the main part of the identification of the asymptotical limits is presented in Section 4. The proof of the convergence part is based on the estimates in [Theorem 2.16, Y2] (in particular the curvature estimate [(2.48), Y2] which is derived from [(7.16), P1]) and the upper bound for l in [Lemma 3.2, Y2].

The identification of the asymptotical limits can be compared with [Theorem 4.9, Y2], which says that a solution of the backward Ricci flow must be a gradient shrinking soliton, provided that the values of its reduced volume at two different times are equal. Recall that, basically, the said equality forces the weakly formulated differential inequality [(2.68), Y2] (i.e. [(7.13), P1]) to become an equality, which then leads to the smoothness of the l -function and implies that the differential inequality [(2.69), Y2] (i.e. [(7.14), P1]) also becomes an equality. These two equalities combined with [Lemma 4.8, Y2] (i.e. [Proposition 9.1, P1]) then lead to the desired gradient shrinking soliton conclusion. We apply this strategy to the current situation. First, in Section 3, we apply Perelman's lower bound for l as presented in [Y2] to show that the asymptotical reduced volume on an asymptotical limit has constant values at all times. Next, in Section 4, we apply the upper bound for Δl and the convergence lemma in Section 2 to obtain the local L^2 strong convergence of ∇l , and use it to derive the differential inequalities [(2.68), Y2] and [(2.69), Y2] in the weak sense for the asymptotical limit. Using this strong convergence we also derive the connection between the two differential inequalities, namely the fact that one becoming an equality forces the other to become an equality. The remaining steps are similar to the proof of [Theorem 4.9, Y2] in [Y2].

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2 Preliminaries

Consider a smooth solution $g = g(\tau)$ of the backward Ricci flow

$$\frac{\partial g}{\partial \tau} = 2Ric \tag{2.1}$$

on an n -dimensional manifold M over an interval $[0, T)$. We assume that $(M, g(\tau))$ is complete for each $\tau \in [0, T)$. Choose an arbitrary point $p \in M$ as the l -base, i.e. the reference point for the l -function, see [Definition 1, Y1].

We'll follow the notations and conventions in [Y1]. In particular, the volume form of a Riemannian metric h on a manifold will be denoted by dq_h or simply dq . The distance function of h will be denoted by d_h or simply d .

Lemma 2.1 *Assume that the curvature operator is nonnegative on $[0, T)$. Let $\bar{\tau} \in (0, T)$. Then there is a positive constant C depending only on the dimension n and the magnitude of $\frac{\bar{\tau}}{T-\bar{\tau}}$, such that the differential inequality $\Delta l \leq C \frac{l}{\tau}$ holds true for each $\tau \in (0, \bar{\tau}]$ in the weak sense, i.e.*

$$\int_M \nabla l \cdot \nabla \phi dq \geq - \int_M C \frac{l}{\tau} \phi dq \quad (2.2)$$

for all nonnegative Lipschitz functions ϕ with compact support, as well as functions ϕ in the Sobolev space $W_{loc}^{1,2}(N)$ with compact support. (Note that C depends only on n if $T = \infty$.)

Proof. By [Theorem 2.20, Y2], the differential inequality

$$\Delta l \leq \frac{|\nabla l|^2}{2} + \frac{R}{2} + \frac{l-n}{2\tau} \quad (2.3)$$

holds true in the weak sense. By the estimates in [Theorem 2.16, Y2], on $[0, \bar{\tau}]$ the right hand side of (2.3) is bounded above by $C \frac{l}{\tau}$. The desired conclusion follows. ■

Next we present an analytic lemma regarding strong convergence of Sobolev functions.

Lemma 2.2 *Let (N, h) be a compact Riemannian manifold, where h denotes the metric. Let f_k be a bounded sequence of functions in the Sobolev space $W^{1,2}(N)$ such that $f_k - f_{k'} \in L^\infty(N)$ for all k, k' , and that*

$$f_k - f_{k'} \rightarrow 0 \quad (2.4)$$

in L^∞ as $k, k' \rightarrow \infty$. Furthermore, assume that there is a positive constant C such that $\Delta f_k \leq C$ for all k in the weak sense, i.e.

$$\int_N \nabla f_k \cdot \nabla \phi dq \geq -C \int \phi dq \quad (2.5)$$

for all nonnegative $\phi \in W^{1,2}(N)$. Then f_k converges strongly in $W^{1,2}(N)$ to a function f . Consequently, we have

$$\int_N |\nabla f_k|^2 \eta dq \rightarrow \int_N |\nabla f|^2 \eta dq \quad (2.6)$$

for each $\eta \in L^\infty(N)$. The same conclusion holds true if $\Delta f_k \leq C$ is replaced by $\Delta f_k \geq -C$.

Proof. Since $f_k \in L^2(N)$, (2.4) implies that f_k is a Cauchy sequence in $L^2(N)$, and hence converges strongly in $L^2(N)$ to a function f . Then the boundedness of f_k in $W^{1,2}(N)$ implies that $f \in W^{1,2}(N)$ and that f_k converges weakly to f . Furthermore, we have $f_k - f \in L^\infty(N)$ and $f_k - f \rightarrow 0$ in $L^\infty(N)$. Taking limit in (2.5) we deduce

$$\int_N \nabla f \cdot \nabla \phi dq \geq -C \int_N \phi dq \quad (2.7)$$

for all nonnegative $\phi \in W^{1,2}(N)$, i.e. $\Delta f \leq C$ in the weak sense.

We set $\delta_k = \sup |f_k - f|$. Then $\delta_k \rightarrow 0$, $f - f_k + \delta_k \geq 0$ and $f_k - f + \delta_k \geq 0$. Choosing $\phi = f - f_k + \delta_k$ in (2.5) we deduce

$$\int_N \nabla f_k \cdot \nabla f dq - \int_N |\nabla f_k|^2 dq \geq -C \int (f - f_k + \delta_k) dq. \quad (2.8)$$

Choosing $\phi = f - f_k + \delta_k$ in (2.7) we deduce

$$\int_N |\nabla f|^2 dq - \int_N \nabla f \cdot \nabla f_k dq \geq -C \int_N (f - f_k + \delta_k) dq. \quad (2.9)$$

It follows that

$$\int_N |\nabla f|^2 dq - \int_N |\nabla f_k|^2 dq \geq 2C \int (f_k - f - \delta_k) dq. \quad (2.10)$$

On the other hand, choosing $\phi = f_k - f + \delta_k$ in (2.5) and (2.7) we infer

$$\int_N |\nabla f|^2 dq - \int_N |\nabla f_k|^2 dq \leq 2C \int (f_k - f + \delta_k) dq. \quad (2.11)$$

Obviously, (2.10) and (2.11) imply

$$\int_N |\nabla f_k|^2 dq \rightarrow \int_N |\nabla f|^2 dq. \quad (2.12)$$

It is easy to see that (2.12) implies that f_k converges strongly to f in $W^{1,2}(N)$. Indeed, we have

$$\int_N |\nabla f_k - \nabla f|^2 dq = \int_N |\nabla f_k|^2 dq + \int_N |\nabla f|^2 dq - 2 \int_N \nabla f_k \cdot \nabla f dq. \quad (2.13)$$

The right hand side is easily seen to converge to zero.

Now we derive (2.6). We have

$$\begin{aligned} \int_N |\nabla f_k|^2 \eta dq - \int_N |\nabla f|^2 \eta dq &= \int_N (\nabla f_k - \nabla f) \cdot \nabla f \eta dq + \\ &\quad \int_N (\nabla f_k - \nabla f) \cdot \nabla f_k \eta dq. \end{aligned} \quad (2.14)$$

Applying Cauchy-Schwarz inequality and the strong convergence of ∇f_k to ∇f in L^2 we then arrive at (2.6).

The proof in the case $\Delta f_k \geq -C$ is similar. ■

It is easy to extend this lemma to the situation of noncompact manifolds. For convenience, we formulate the result in the set-up of pointed manifolds.

Lemma 2.3 *Let (N, h, p) be a pointed complete noncompact Riemannian manifold. Let f_k be a locally bounded sequence of functions in $W_{loc}^{1,2}(N)$ such that the Δf_k are locally uniformly bounded from above in the weak sense, namely there exist for each $A > 0$ positive constants $C_1(A)$ and $C_2(A)$ such that for each k the $W^{1,2}$ -norm of f_k on the geodesic ball $d(p, \cdot) < A$ is bounded above by $C_1(A)$, and $\Delta f_k \leq C_A$ holds true on $d(p, \cdot) < A$ in the weak sense, i.e.*

$$\int_{N_k} \nabla f_k \cdot \nabla \phi dq \geq -C_A \int_{N_k} \phi dq \quad (2.15)$$

for all nonnegative $\phi \in W_{loc}^{1,2}(N_k)$ with support contained in $d(p, \cdot) < A$. Furthermore, assume that $f_k - f_{k'} \in L_{loc}^\infty(N)$ for all k, k' , and $f_k - f_{k'} \rightarrow 0$ in $L_{loc}^\infty(N)$. Then f_k converges strongly in $W_{loc}^{1,2}$ to a function f . Consequently, we have

$$\int_N |\nabla f_k|^2 \eta dq \rightarrow \int_N |\nabla f|^2 \eta dq \quad (2.16)$$

for each $\eta \in L^\infty(N)$ with compact support.

The same conclusion holds true if the Δf_k are assumed to be locally uniformly bounded from below in the weak sense instead of from above.

Proof. First, f_k converges to some f weakly in $W_{loc}^{1,2}(N)$. Moreover, $f_k - f$ converges to zero in $L_{loc}^\infty(N)$. As in the proof of Lemma 2.2 we obtain by taking limit in (2.15)

$$\int_N \nabla f \cdot \nabla \phi dq \geq -C \int_N \phi dq \quad (2.17)$$

for all nonnegative $\phi \in W_{loc}^{1,2}(N)$ with compact support. For a fixed $A > 0$, we set $\delta_k = \sup_{d(p, \cdot) \leq 2A} |f_k - f|$. Let ψ be a nonnegative Lipschitz function on N which takes the value 1 on $d(p, \cdot) \leq A$ and is zero on $d(p, \cdot) \geq 2A$. Choosing $\phi = \psi(f - f_k + \delta_k)$ in (2.15) and (2.17) we deduce

$$\begin{aligned} \int_N |\nabla f|^2 dq - \int_N |\nabla f_k|^2 dq &\geq 2C_A \int \psi(f_k - f - \delta_k) dq + \int_N \nabla f_k \cdot \nabla \psi(f_k - f - \delta_k) dq \\ &\quad + \int_N \nabla f \cdot \nabla \psi(f_k - f - \delta_k) dq. \end{aligned} \quad (2.18)$$

On the other hand, choosing $\phi = \psi(f_k - f + \delta_k)$ in (2.15) and (2.17) we infer

$$\begin{aligned} \int_N |\nabla f|^2 \psi dq - \int_N |\nabla f_k|^2 \psi dq &\leq 2C_A \int_N \psi(f - f_k + \delta_k) dq \\ + \int_N \nabla f_k \cdot \nabla \psi(f_k - f + \delta_k) dq &+ \int_N \nabla f \cdot \nabla \psi(f_k - f + \delta_k) dq. \end{aligned} \quad (2.19)$$

We can use the local boundedness of f_k in $W_{loc}^{1,2}$ and Cauchy-Schwarz inequality to handle the second and third terms on the right hand side of (2.18) and (2.19). It

follows that $\int_N |\nabla f_k|^2 \psi dq \rightarrow \int_N |\nabla f|^2 \psi dq$. Then

$$\begin{aligned} \int_{d(p,\cdot) \leq A} |\nabla f_k - \nabla f|^2 dq &\leq \int_N |\nabla f_k - \nabla f|^2 \psi dq = \int_N |\nabla f_k|^2 \psi dq + \int_N |\nabla f|^2 \psi dq \\ &\quad - 2 \int_N \nabla f_k \cdot \nabla f \psi dq \rightarrow 0. \end{aligned} \quad (2.20)$$

The derivation of (2.16) is similar to that of (2.6) in the proof of Lemma 2.2.

The case that Δf_k are locally uniformly bounded from below in the weak sense is similar. \blacksquare

Next we extend the above results to the situation of a sequence of converging Riemannian manifolds. For simplicity of formulation we state a weaker version of the result.

Lemma 2.4 *Let (N_k, h_k, p_k) be a sequence of pointed complete Riemannian manifolds converging in C^1 to a pointed complete Riemannian manifold (N, h, p) . Let $f_k \in W_{loc}^{1,2}(N_k) \cap L_{loc}^\infty(N_k)$ which converges in L_{loc}^∞ and weakly in $W_{loc}^{1,2}$ to a function $f \in W_{loc}^{1,2}(N) \cap L_{loc}^\infty(N)$. Furthermore, we assume that the Δf_k are locally uniformly bounded from above in the weak sense or locally uniformly bounded from below in the weak sense. Then f_k converges strongly in $W_{loc}^{1,2}$ to f . Consequently, we have*

$$\int_{N_k} |\nabla f_k|^2 \eta_k dq \rightarrow \int_N |\nabla f|^2 \eta dq, \quad (2.21)$$

whenever $\eta \in L^\infty(N)$ and $\eta_k \in L^\infty(N_k)$ with support contained in $d_{h_k}(p_k, \cdot) < A$ for a fixed $A > 0$, such that η_k converges to η in L^∞ .

Proof. The proof of Lemma 2.3 can easily be carried over. Note that the pointed convergence of (N_k, h_k, p_k) to (N, h, p) involves suitable embeddings from increasing domains of N_k into N . Using these embeddings the functions f_k can be transplanted to increasing domains of N . Then the difference between f_k and f can be measured. \blacksquare

3 Asymptotic Limits of κ -Solutions I

Let $\tilde{g}(t)$, $-\infty < t \leq 0$ be a κ -solution for some $\kappa > 0$ on a manifold M . Recall [P1] that this means that $\tilde{g}(t)$ is a smooth solution of the Ricci flow

$$\frac{\partial g}{\partial t} = -2Ric \quad (3.1)$$

on $M \times (-\infty, 0]$, such that for each t , the metric $\tilde{g}(t)$ is a complete non-flat metric of bounded and nonnegative curvature operator. Moreover, each $\tilde{g}(t)$ is κ -noncollapsed

on all scales. (See [P1] for the definition of κ -noncollapsedness.) To understand the structures of $\tilde{g}(t)$, we analyse its rescaled asymptotical limits at the time infinity. One needs to use a blow-down rescaling, because the Ricci flow equation and nonnegative curvature imply that $\tilde{g}(t)$ expands as t decreases.

Pick an arbitrary time $t_0 \leq 0$. Then $g(\tau) = \tilde{g}(t_0 - \tau)$, $\tau \in [0, \infty)$ is a solution of the backward Ricci flow (2.1) on $M \times [0, \infty)$. Obviously, g retains the properties of \tilde{g} , namely g is κ -noncollapsed on all scales and has bounded and nonnegative curvature operator.

Next we choose an arbitrary $p \in M$ as the l -base, i.e. the reference point for the l -function l and the reduced volume \tilde{V} (see [Definition 1, Y2]). We set

$$\tilde{V}^\infty = \lim_{\tau \rightarrow \infty} \tilde{V}(\tau). \quad (3.2)$$

By [Theorem 4.5, Y2], this limit exists and is finite.

For $a > 0$, we set $g_a(\tau) = \frac{1}{a}g(a\tau)$. Choose a base point $\bar{q} \in M$ we then have a pointed solution of the backward Ricci flow $(g_a, M \times (0, \infty), \bar{q})$. The following theorem is precisely [Proposition 11.2, P1].

Theorem 3.1 *Let $\tau_k \rightarrow \infty$ be given. For each τ_k , let $x(\tau_k)$ be a minimum point of $l(\cdot, \tau)$. Then the pointed flows $(g_{\tau_k}, M \times (0, \infty), x(\tau_k))$ subconverge smoothly to pointed nonflat gradient shrinking solitons $(g_\infty, M_\infty \times (0, \infty), x_\infty)$.*

In the remainder of this section, we establish the convergence part of Theorem 3.1, and obtain some basic properties for the limit solutions. For convenience, we formulate the convergence part as a proposition.

Proposition 3.2 *The pointed flows $(g_{\tau_k}, M \times (0, \infty), x(\tau_k))$ subconverge smoothly to pointed solutions $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ of the backward Ricci flow. The limits will be called asymptotical limits of \tilde{g} or g .*

Proof. The various quantities associated with g_{τ_k} will be indicated by the subscript k or g_{τ_k} , e.g. $l_k = l_{g_{\tau_k}}$, $\tilde{V}_k = \tilde{V}_{g_{\tau_k}}$ (the l -function and reduced volume associated with g_{τ_k}), and $d_k = d_{g_{\tau_k}}$.

First note that by [Lemma 2.3, Y2] or [Lemma 3.2, Y2] the minimum points $x(\tau_k)$ exist. By [Lemma 3.1, Y2] and the scaling invariance of the l -function we have

$$l_k(x(\tau_k), 1) \leq \frac{n}{2}. \quad (3.3)$$

By this estimate and [Lemma 3.2, Y2] we infer

$$l_k(q, 1) \leq C_2 d_k^2(x(\tau_k), q, 1) + n \quad (3.4)$$

for all $q \in M$, where C_2 is a positive constant depending only on the dimension n . Then it follows from the Harnack inequality [(2.52), Y2] that

$$l_k(q, \tau) \leq \tau^{\pm C} (C_2(d_k^2(x(\tau_k), q, 1) + n), \quad (3.5)$$

where $\pm = +$ if $\tau \geq 1$, $\pm = -$ if $\tau < 1$, and C is a positive constant depending only on n . Consequently, we obtain from [Theorem 2.16, Y2] and the nonnegativity of curvature operator the estimate

$$|Rm|_k(q, \tau) \leq C\tau^{-1 \pm C} (d_k^2(x(\tau_k), q, 1) + 1). \quad (3.6)$$

By the κ -noncollapsing assumption and (3.6), we obtain for each k and $\tau \in (0, \infty)$ local injectivity radius estimates for $g_{\tau_k}(\tau)$ which depend only on $d_k(x(\tau_k), \cdot, 1)$, τ (in terms of a positive lower bound and an upper bound of τ), κ and n . It is then straightforward to apply Gromov-Cheeger-Hamilton compactness [H] to obtain pointed smooth subconvergence of $(g_{\tau_k}, M \times (0, \infty), x(\tau_k))$ to pointed solutions $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ of the backward Ricci flow. \blacksquare

Let $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ be an asymptotical limit. The corresponding converging subsequence will still be denoted by g_{τ_k} . By [Theorem 2.16, Y2], the rescaling invariance and (3.5), the l_k subconverge locally uniformly to locally Lipschitz functions l_∞ on $M_\infty \times (0, \infty)$.

Definition 1 The limits l_∞ will be called *limit l -functions*.

Lemma 3.3 *Let $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ be an asymptotical limit and l_∞ a limit l -function. Then $l_\infty(\cdot, \tau)$ is locally Lipschitz for every τ , and $l_\infty(q, \cdot)$ is locally Lipschitz for every q . Moreover, we have*

$$R_\infty \leq C \frac{l_\infty}{\tau} \quad (3.7)$$

everywhere,

$$|\nabla l_\infty| \leq C \frac{l_\infty}{\tau} \quad (3.8)$$

almost everywhere in M_∞ for each $\tau \in (0, \infty)$ (the ∇ and the norm are those of g_∞),

$$\left| \frac{\partial l_\infty}{\partial \tau} \right| \leq C \frac{l_\infty}{\tau} \quad (3.9)$$

almost everywhere in $(0, \infty)$ for each $q \in M_\infty$, and

$$\left(\frac{\tau_1}{\tau_2} \right)^C \leq \frac{l_\infty(q, \tau_2)}{l_\infty(q, \tau_1)} \leq \left(\frac{\tau_2}{\tau_1} \right)^C \quad (3.10)$$

for all $q \in M$ and $0 < \tau_1 < \tau_2$, where $R_\infty = R_{g_\infty}$ (the scalar curvature of g_∞) and C depends only on n .

Proof. These follow from the local uniform convergence of l_k to l_∞ and the corresponding properties of the l -function as given in [Theorem 2.16, Y2]. The estimates for the gradient and the τ -derivative are obtained in terms of estimating the relevant Lipschitz constants. \blacksquare

Lemma 3.4 *We have*

$$C_1 \frac{d_k^2(x(\tau_k), q, \tau)}{\tau} - \frac{n}{2} \tau^{\pm C} - 1 \leq l_k(q, \tau) \leq C_2 \frac{d_k^2(x(\tau_k), q, \tau)}{\tau} + n \tau^{\pm C} \quad (3.11)$$

for all $q \in M, \tau \in (0, \infty)$ and

$$C_1 \frac{d_\infty^2(x_\infty, q, \tau)}{\tau} - \frac{n}{2} \tau^{\pm C} - 1 \leq l_\infty(q, \tau) \leq C_2 \frac{d_\infty^2(x_\infty, q, \tau)}{\tau} + n \tau^{\pm C} \quad (3.12)$$

for all $q \in M_\infty, \tau \in (0, \infty)$, where d_∞ denotes the distance function of g_∞ , and the positive constant C_1, C_2 and C depend only on the dimension n .

Proof. By the Harnack inequality [(2.52), Y2] and (3.3) we deduce

$$l_k(x(\tau_k), \tau) \leq \frac{n}{2} \tau^{\pm C}. \quad (3.13)$$

Applying [Lemma 3.2, Y2] we then arrive at (3.11). Taking limit leads to (3.12). (Note that a similar estimate was employed in the proof of [Theorem 3.3, Y2] in [Y2].) \blacksquare

Henceforth, we'll fix a limit l -function l_∞ for each asymptotical limit. Obviously, the conclusions we derive are valid for each choice of the limit l -function.

Definition 2 Let $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ be an asymptotical limit together with a limit l -function l_∞ . We define the *asymptotical reduced volume* to be

$$\tilde{V}_\infty(\tau) = \int_{M_\infty} e^{-l_\infty} \tau^{-\frac{n}{2}} dq, \quad (3.14)$$

where $dq = dq_{g_\infty}$.

Lemma 3.5 $\tilde{V}_k(\tau)$ converges to $\tilde{V}_\infty(\tau)$ for each $\tau > 0$ and hence there holds

$$\tilde{V}_\infty(\tau) = \tilde{V}^\infty \quad (3.15)$$

for each τ .

Proof. Fix $\tau > 0$. For each $A > 0$ we have

$$\tilde{V}_k(\tau) = \int_{d_k(x(\tau_k), \cdot, \tau) \leq A} e^{-l_k \tau^{-\frac{n}{2}}} dq_k + \int_{d_k(x(\tau_k), \cdot, \tau) > A} e^{-l_k \tau^{-\frac{n}{2}}} dq_k \quad (3.16)$$

and

$$\tilde{V}_\infty(\tau) = \int_{d_\infty(x_\infty, \cdot, \tau) \leq A} e^{-l_\infty \tau^{-\frac{n}{2}}} dq + \int_{d_\infty(x_\infty, \cdot, \tau) > A} e^{-l_\infty \tau^{-\frac{n}{2}}} dq. \quad (3.17)$$

As in the proof of [Theorem 3.3, Y2] in [Y2], the lower bound in (3.11) and volume comparison imply

$$\int_{d_k(x(\tau_k), \cdot, \tau) > A} e^{-l_k \tau^{-\frac{n}{2}}} dq_k \leq C e^{-cA} \quad (3.18)$$

for suitable positive constants C and c depending only on n . Similarly, the lower bound in (3.12) and volume comparison imply

$$\int_{d_\infty(x_\infty, \cdot, \tau) > A} e^{-l_\infty \tau^{-\frac{n}{2}}} dq \leq C e^{-cA}. \quad (3.19)$$

(We can choose the constants here to be the same as in (3.18).) By the pointed convergence of $(g_{\tau_k}, M \times (0, \infty), x(\tau_k))$ to $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ and the convergence of l_k to l_∞ we then infer $\lim | \tilde{V}_k(\tau) - \tilde{V}_\infty(\tau) | \leq 2C e^{-cA}$. Since A is arbitrary, we conclude that $\lim \tilde{V}_k(\tau) = \tilde{V}_\infty(\tau)$. But $\tilde{V}_k(\tau) = \tilde{V}(\tau_k \tau) \rightarrow \tilde{V}^\infty$. Hence we arrive at (3.15). \blacksquare

4 Asymptotical Limits of κ -solutions II

The goal of this section is to establish the soliton characterization part of Theorem 3.1. We formulate it as a proposition. See [Y2] for the terminologies regarding gradient shrinking solitons used below.

Proposition 4.1 *Each asymptotical limit $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ is a pointed nonflat gradient shrinking soliton with time origin 0. Moreover, the limit l -functions l_∞ are potential functions.*

To prove this proposition, we need a few lemmas. We fix an asymptotical limit $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ together with a limit l -function l_∞ . We have the corresponding converging subsequences $(g_{\tau_k}, M \times (0, \infty), x(\tau_k))$ and l_k .

Lemma 4.2 *For each $\tau > 0$, $l_k(\cdot, \tau)$ converges strongly to $l_\infty(\cdot, \tau)$ in $W_{loc}^{1,2}$. Consequently, we have*

$$\int_{M, g_{\tau_k}} |\nabla l_k|^2 \phi_k dq|_\tau \rightarrow \int_{M_\infty} |\nabla l_\infty|^2 \phi dq, \quad (4.1)$$

whenever $\phi \in L^\infty(M_\infty)$ and $\phi_k \in L^\infty(M)$ with support contained in $d_k(x(\tau_k), \cdot, \tau) < A$ for a fixed $A > 0$, such that ϕ_k converges in L^∞ to ϕ .

Proof. By [Theorem 2.16, Y2] and (3.5), l_k is a locally bounded sequence in $W_{loc}^{1,2}$. Since it converges to l_∞ in L_{loc}^∞ , it also converges to l_∞ weakly in $W_{loc}^{1,2}$. Hence the conclusions of the lemma follow from Lemma 2.1 and Lemma 2.4. \blacksquare

We also need the L_{loc}^2 strong convergence of ∇l_k over the spacetime $M \times (0, \infty)$.

Lemma 4.3 *The vector fields ∇l_k on $M \times (0, \infty)$ converge to the vector field ∇l_∞ on $M_\infty \times (0, \infty)$ strongly in L_{loc}^2 . Consequently, we have for arbitrary $\tau_2 > \tau_1 > 0$*

$$\int_{\tau_1}^{\tau_2} \int_{M, g_{\tau_k}} |\nabla l_k|^2 \phi_k dq d\tau \rightarrow \int_{\tau_1}^{\tau_2} \int_{M_\infty} |\nabla l_\infty|^2 \phi dq d\tau, \quad (4.2)$$

whenever $\phi \in L^\infty(M_\infty \times [\tau_1, \tau_2])$ and $\phi_k \in L^\infty(M \times [\tau_1, \tau_2])$ with support contained in $\{d_k(x(\tau_k), \cdot, 1) < A\} \times [\tau_1, \tau_2]$ for a fixed $A > 0$, such that ϕ_k converges to ϕ in L^∞ .

Proof. By [Theorem 2.16, Y2] and (3.5), l_k is a locally bounded sequence in $W_{loc}^{1,2}$ over the spacetime. It follows that l_k converges weakly in $W_{loc}^{1,2}$ over the spacetime. Now we can argue as in the proof of Lemma 2.3, employing integrations over the spacetime with respect to the volume form $dq d\tau$. Note that the involved Δ and ∇ are along the space with a time dependence on τ . \blacksquare

Lemma 4.4 *The equation*

$$\frac{\partial l_\infty}{\partial \tau} - \frac{R_\infty}{2} + \frac{|\nabla l_\infty|^2}{2} + \frac{l_\infty}{2\tau} = 0 \quad (4.3)$$

holds true almost everywhere on $M_\infty \times (0, \infty)$. The inequality

$$\Delta l_\infty - \frac{|\nabla l_\infty|^2}{2} + \frac{R_\infty}{2} + \frac{l_\infty - n}{2\tau} \leq 0 \quad (4.4)$$

holds true for each $\tau > 0$ in the weak sense, i.e.

$$\int_{M_\infty} \left\{ -\nabla l_\infty \cdot \nabla \phi + \frac{1}{2}(-|\nabla l_\infty|^2 + R_\infty + \frac{l_\infty - n}{\tau})\phi \right\} dq \leq 0 \quad (4.5)$$

for all nonnegative Lipschitz functions ϕ with compact support. Finally, the inequality

$$\frac{\partial l_\infty}{\partial \tau} - \Delta l_\infty + |\nabla l_\infty|^2 - R_\infty + \frac{n}{2\tau} \geq 0 \quad (4.6)$$

holds true on $M_\infty \times (0, \infty)$ when Δ is interpreted in the weak sense, i.e.

$$Q_{\tau_1, \tau_2}(\phi) \geq 0 \quad (4.7)$$

for arbitrary $\tau_2 > \tau_1 > 0$ and nonnegative Lipschitz functions ϕ on $M_\infty \times [\tau_1, \tau_2]$ with compact support, where

$$Q_{\tau_1, \tau_2}(\phi) = \int_{\tau_1}^{\tau_2} \int_{M_\infty} \left\{ \nabla l_\infty \cdot \nabla \phi + \left(\frac{\partial l_\infty}{\partial \tau} + |\nabla l_\infty|^2 - R_\infty + \frac{n}{2\tau} \right) \phi \right\} dq d\tau. \quad (4.8)$$

Proof. First consider a fixed $\tau > 0$. Let ϕ be a nonnegative smooth function on M_∞ with compact support. Let ϕ_k be a sequence of nonnegative smooth functions on M such that the support of ϕ_k is contained in $d_k(x(\tau_k), \cdot, \tau) < A$ for a fixed $A > 0$, and that ϕ_k converges smoothly to ϕ . By [Theorem 2.20, Y2] we have for each k

$$\int_{M, g_{\tau_k}} \left\{ -\nabla l_k \cdot \nabla \phi_k + \left(-\frac{|\nabla l_k|^2}{2} + \frac{R_k}{2} + \frac{l_k - n}{2\tau} \right) \phi_k \right\} dq \leq 0 \quad (4.9)$$

On account of Lemma 4.2 and the pointed convergence of $(g_{\tau_k}, M \times (0, \infty), x(\tau_k))$ to $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ we can take limit in (4.9) to arrive at (4.5) with the above ϕ . Then it follows that (4.5) also holds true for nonnegative Lipschitz or Sobolev functions ϕ with compact support.

Next consider $\tau_2 > \tau_1 > 0$ and a smooth nonnegative function ϕ on $M_\infty \times [\tau_1, \tau_2]$ with compact support. Let ϕ_k be a sequence of smooth functions on $M \times [\tau_1, \tau_2]$ such that the support of ϕ_k is contained in $d_k(x(\tau_k), \cdot, 1) < A$ for a fixed $A > 0$, and that ϕ_k converges smoothly to ϕ . By [Theorem 2.20, Y2] we have

$$\int_{\tau_1}^{\tau_2} \int_{M, g_{\tau_k}} \left\{ \nabla l_k \cdot \nabla \phi_k + \left(\frac{\partial l_k}{\partial \tau} + |\nabla l_k|^2 - R_k + \frac{n}{2\tau} \right) \phi_k \right\} dq d\tau \geq 0. \quad (4.10)$$

Since $\frac{\partial}{\partial \tau} dq = R_k dq$ we can rewrite this equation as follows

$$\begin{aligned} & \int_{M, g_{\tau_k}} l_k \phi_k dq \Big|_{\tau_2} - \int_{M, g_{\tau_k}} l_k \phi_k dq \Big|_{\tau_1} + \\ & \int_{\tau_1}^{\tau_2} \int_{M, g_{\tau_k}} \left\{ \nabla l_k \cdot \nabla \phi_k + \left(-l_k \frac{\partial \phi_k}{\partial \tau} + |\nabla l_k|^2 - 2R_k + \frac{n}{2\tau} \right) \phi_k \right\} dq d\tau \geq 0. \end{aligned} \quad (4.11)$$

On account of Lemma 4.3, the local uniform convergence of l_k to l_∞ and the pointed convergence of $(g_{\tau_k}, M \times (0, \infty), x(\tau_k))$ we can take limit in (4.11) to arrive at

$$\begin{aligned} & \int_{M_\infty} l_\infty \phi dq \Big|_{\tau_2} - \int_{M_\infty} l_\infty \phi dq \Big|_{\tau_1} + \\ & \int_{\tau_1}^{\tau_2} \int_{M_\infty} \left\{ \nabla l_\infty \cdot \nabla \phi + \left(-l_\infty \frac{\partial \phi}{\partial \tau} + |\nabla l_\infty|^2 - 2R_\infty + \frac{n}{2\tau} \right) \phi \right\} dq d\tau \geq 0, \end{aligned} \quad (4.12)$$

which is equivalent to (4.7) for the above ϕ . Then (4.7) also holds true for nonnegative Lipschitz or Sobolev ϕ with compact support.

By [Lemma 2.19, Y2] and [Lemma 2.15, Y2] we have

$$\frac{\partial l_k}{\partial \tau} - \frac{R_k}{2} + \frac{|\nabla l_k|^2}{2} + \frac{l_k}{2\tau} = 0 \quad (4.13)$$

almost everywhere on $M \times (0, \infty)$. We turn this equation into an integral equation in the same fashion as (4.11). Taking limit we then arrive at a corresponding integral

equation for l_∞ , which implies that (4.3) holds true almost everywhere. \blacksquare

Proof of Proposition 4.1

Fix an asymptotical limit $(g_\infty, M_\infty \times (0, \infty), x_\infty)$ together with a limit l -function l_∞ as above. By (4.7), Lemma 3.4 and the arguments in the proof of [Theorem 3.3, Y2] in [Y2] we infer that $Q_{\tau_1, \tau_2}(\phi) \geq 0$ holds true for nonnegative admissible locally Lipschitz ϕ on $M_\infty \times [\tau_1, \tau_2]$, where ‘‘admissible’’ means that $\phi \leq C\tau^{-\frac{n}{2}}e^{-l_\infty}$ and $|\nabla\phi| \leq C\tau^{-\frac{n+1}{2}}\sqrt{l_\infty}e^{-l_\infty}$ for all $q \in M_\infty$ and $\tau \in [\tau_1, \tau_2]$ and some positive constant C depending on ϕ . On the other hand, we can argue as in the proof of [Lemma 4.6, Y2] in [Y2] to deduce

$$\tilde{V}_\infty(\tau_2) - \tilde{V}_\infty(\tau_1) = - \int_{\tau_1}^{\tau_2} \int_{M_\infty} \left(\frac{\partial l_\infty}{\partial \tau} - R_\infty + \frac{n}{2\tau} \right) e^{-l_\infty} \tau^{-\frac{n}{2}} dq d\tau. \quad (4.14)$$

By Lemma 3.5 we then infer that $Q(\tau^{-\frac{n}{2}}e^{-l_\infty}) = 0$. As in the proof of [Theorem 4.9, Y2] in [Y2], we then deduce that $Q_{\tau_1, \tau_2}(\phi) = 0$ for all admissible locally Lipschitz ϕ . By basic regularity theory for parabolic equations we derive that l_∞ is smooth and the equation

$$\frac{\partial l_\infty}{\partial \tau} - \Delta l_\infty + |\nabla l_\infty|^2 - R_\infty + \frac{n}{2\tau} = 0 \quad (4.15)$$

holds true everywhere. By Lemma 4.4, the equation (4.3) also holds true everywhere. Taking the difference between (4.3) and (4.15) we then arrive at the equation

$$\Delta l_\infty - \frac{|\nabla l_\infty|^2}{2} + \frac{R_\infty}{2} + \frac{l_\infty - n}{2\tau} = 0. \quad (4.16)$$

By (4.15), (4.16) and [Lemma 4.8, Y2] we finally conclude that g_∞ is a gradient shrinking soliton with time origin 0 and potential function l_∞ .

Now we show that g_∞ is nonflat. Assume that g_∞ is flat. Then the soliton equation [(4.16), Y2] implies

$$\nabla_{g_\infty}^2 l_\infty = \frac{1}{2\tau} g_\infty. \quad (4.17)$$

It follows that $l_\infty(\cdot, 1)$ is a strictly convex function with a unique minimum point (by the construction it has a minimum point x_∞). By Lemma 3.4, its level sets are compact. Employing the gradient flow of $l_\infty(\cdot, 1)$ one readily shows that M_∞ is diffeomorphic to \mathbf{R}^n and then each $(M_\infty, g_\infty(\tau))$ is isometric to \mathbf{R}^n . Next we observe that (4.17) implies that $\Delta_{g_\infty} l_\infty = \frac{n}{2\tau}$, which together with (4.16) leads to

$$|\nabla l_\infty|^2 = \frac{l_\infty}{\tau}. \quad (4.18)$$

For a fixed τ , we identify (M_∞, g_∞) with \mathbf{R}^n via an isometry. Let x denote the coordinates on \mathbf{R}^n . Then (4.17) implies that

$$\nabla l_\infty(x, \tau) = \nabla \frac{|x|^2}{4\tau} + v(\tau) \quad (4.19)$$

for a constant vector $v(\tau)$. Hence

$$l_\infty(x, \tau) = \frac{|x|^2}{4\tau} + v(\tau) \cdot x + c(\tau) \quad (4.20)$$

for a constant $c(\tau)$. A simple calculation using (4.18), (4.19) and (4.20) yields $c(\tau) = \tau|v|^2$. It follows that $l_\infty(x, \tau) = \frac{|x+2\tau v|^2}{4\tau}$ and hence $\tilde{V}_\infty(\tau) = (4\pi)^{\frac{n}{2}}$ for each τ . But $\tilde{V}_\infty(\tau) = \tilde{V}^\infty < (4\pi)^{\frac{n}{2}}$ by [Theorem 4.4, Y2] and [Theorem 4.5, Y2]. Thus we arrive at a contradiction. ■

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