Lecture 17 <u>Announcements</u>: • Revision due by March 4th Recall:

Del: I Weak solution of continuity eqn): Given $v: \mathbb{R}^{d} \times (0,T) \to \mathbb{R}^{d}$ meas, $u \in \mathcal{M}(\mathbb{R}^{d})$, $u:(0,T) \to \mathcal{M}(\mathbb{R}^{d})$ is a weak solution of the continuity equation $\begin{cases} \partial_{\pm}\mu + \nabla \cdot (\mu v) = 0 \\ \mu(0) = \mu_0 \end{cases}$

· ult) is narrowly cts, lime ult)= uo narrowly · the PDE holds in distribution, that is $\int \int (\partial_t \theta(x,t) + \nabla \theta(x,t) \cdot v(x,t)) d\mu_t(x) dt = 0$ $\forall \mathcal{Q} \in C^{\infty}_{c}(\mathbb{R}^{d} \times (0,T)).$

Def: (solution of ODE): Given $\nabla (\mathbb{R}^{d} \times (0,T) \rightarrow \mathbb{R}^{d} \text{ meas and } \tau_{o} \in \mathbb{R}^{d},$ $\chi: (0,T) \rightarrow \mathbb{R}^{d}$ is a solution of

$$\begin{cases} \chi'/t \ = \sqrt{\chi/t}, t \\ \chi(0) = \chi_0 \end{cases}$$

if
•
$$\chi(t)$$
 is locally abscts, $\frac{1}{2}S_0 + \chi(t) = \chi_0$.
• the ODE holds in the integral sense, i.e.
 $\chi(t) = \chi_0 + \int_{0}^{t} \chi(\chi(s), s) ds$, $\forall t \in [0,T]$.

Fix
$$v(x,t)$$
. Suppose that a solution of
(ODE) exists $\forall x_0 \in \mathbb{R}^d$. In this
Case, we may consider the flow map
induced by v_{i} ,
 $\chi_{v(y)} = \chi(y,t) = \chi(t)$, where $\chi(t)$ is solv $w/\chi(0) = y$.

Prop: Fix $\tau(x,t)$ and $\mu_0 \in P_2(\mathbb{R}^d)$. Suppose that solves of (ODE) exist μ_0 -a.e. $x_0 \in \mathbb{R}^d$, $\forall t \in [0,T]$. Let χ_t be the corresponding flow map, which is defined μ_0 -a.e.

Then define $\mu_t = \chi_t \# \mu_0$. Then, if $SS_{Re}|v(x,t)|^2 d\mu(x) dt < +\infty$, • $\mu \in P_2(\mathbb{R}^d) \quad \forall \quad t \in [0,T]$ · (m,v) is a soln of ((E) on TRd×[0,T]. Pg: Last time, we showed second bullet. Furthermore, $\mu_t \in P(\mathbb{R}^d) \ \forall \ t \in [0,T].$ To see that $S[x]^2 d\mu_t(x) < +\infty$ $\forall t \in [0,T],$ recall that, by defn of ODE, |Xtly)|= |y| + Shr(Xsly), S)ps no-a.e. y. $S|\chi|^2 d\mu_t(\chi) = S|\chi_t(\chi)|^2 d\mu_0(\chi)$ = 2S[y]^2 d\mu_0(\chi) +2CT S[v(\chi_s(\chi),s) bd_{\mu_0}(\chi) ds R^a

= $2 \int \left[y \left[\frac{d}{d} y \right] + 2 C_T \int \int \left[\frac{v(x_s)}{R^2} \right]^2 d\mu_t(x) ds$ $<+\infty$

We will now use this Lagrangian perspective to connect Wasserstein geodesics to (CE). Def: Given a metric space (χ, d) , $\chi: (0, 1) \rightarrow \chi$ is a (constant speed) <u>geodesic</u> if d(x(t), x(s)) = |t-s|d(x(0), x(1))|Brop: Given Mo, M, EP2(Rd), Mo, M1 ~ Jd, there exists a constant speed geodesic between them. Furthermore, there exists a velocity v s.t. (μ, v) is a solu (CE) and $(SIv(x,t))^{2}d\mu_{t}(x))^{2} \equiv W_{2}(\mu_{0},\mu_{1}) \forall t \in [0,1]$

The proof of this proposition relies on the following: Prop: Suppose mv EP2(Re) are abscts w.r.t. Lebesque. Let T denote the OT map from u to V (for the quadratic cost). (i) $\delta = (id \times T) \# \mu$ is the unique OT plan. (ii) If T is the OT map from $\nu + \sigma \mu$, to T=id vale, ToT=id prale. In particular, T is left invertible mate. with $T^{-1} = \tilde{T}$. PR: First show (i). We have already shown & is an OT plan. It remains to show it is unique. By Brenier's theorem, J 9eL 1/w) proper, convex, lsc s.t. T= V9. Thus by defn, $\chi^2 = \nabla \mathcal{Q}(\chi^2) \in \partial \mathcal{Q}(\chi^2)$ X-a.e. Fix another OT plan y'.

By Knott-Smith theorem,

$$x^2 \in \mathcal{SP}(x^1)$$
 $\mathcal{S}'-a.e.$
Since \mathcal{P} is diff max.
 $x^2 = \nabla \mathcal{P}(x^1)$ $\mathcal{S}'-a.e.$
Thus $\mathcal{S}'=(id \times T) \# u = \mathcal{S}.$
clow, show (ii).
Since $\mathcal{S}'=(\mathcal{T} \times id) \# \mathcal{V}$ is an OT
plan, by part (i),
 $\mathcal{S}'=(\mathcal{T} \times id) \# \mathcal{V}=(id \times T) \# u = \mathcal{S}.$
Thus,
 $\mathcal{S}'=(\mathcal{T} \times id) \# \mathcal{V}=(id \times T) \# u = \mathcal{S}.$

Thus,

$$Sly = T \cdot T(y) dv(y)$$

 $= Slx^2 - T(x^2) l dv'(x^2, x^2)$
 $= Slx^2 - T(x^2) l dv(x^2, x^2)$
 $= S \cdot T(y) - T(y) l dv(y) = 0.$

Arguing Symmetrically then gives the result. Now, we prove main geoclesic thm. Since M, u, << fd, J OT map T from No to M1. idha)=x idk.)=x Define $u(t) := ((1-t)id + tT) # u_0$ Note that, if $T = \nabla P$, $T_t = \nabla ((1-t)\frac{x}{2} + tP(x))$. Thus, by change of variables formula, $d\mu_t(y) = \frac{\mu_0}{|\det DT_t|} \circ T_t'(y) d\mathcal{I}^q(y) |_{T_t(\mathbb{R}^d)}$ Thus ME << I for all teloit.

Since $T_t # \mu_0 = \mu_t \quad \forall t \in [0, \overline{D}]$ $(T_t \times T_s) # \mu_0 \in \Gamma(\mu_t, \mu_s)$. Thus, $W_{2}^{2}(\mu_{t},\mu_{s}) \leq \int |\chi^{1} - \chi^{2}|^{2} d(T_{t} \times T_{s}) \# \mu_{0}(\chi^{2},\chi^{2})$ $= S[T_t(x) - T_s(x)]^2 d\mu_o(x)$ $= |t-s|^2 \int |x-T(x)|^2 d\mu_0(x)$ $= |t-s|^2 W_2^2(\mu_0,\mu_1)$ For the other direction of the inequality, note that, 0+25 $W_2(\mu_0,\mu_1) \stackrel{\scriptstyle{\scriptstyle{\leftarrow}}}{=} W_2(\mu_0,\mu_s) \stackrel{\scriptstyle{\scriptstyle{\leftarrow}}}{=} W_2(\mu_0,\mu_s) \stackrel{\scriptstyle{\scriptstyle{\leftarrow}}}{=} W_2(\mu_0,\mu_s) \stackrel{\scriptstyle{\scriptstyle{\leftarrow}}}{=} W_2(\mu_0,\mu_s)$ $\leq [s + (t - s) + (1 - t)] W_2(\mu_0, \mu_1)$ = W2(μ_0, μ_1) we must have "=" throughout Thus,

Thus, between any u.u., a constant speed geodesic exists. To now show that this constant speed geodesic solves (CE), recall from previous prop, T_t is left invertible μ_0 -a.e. and T_t^{-1} is of map from μ_t to μ_0 .

Define,
$$v(x,t) = T \circ T_t^{-1}(x) - T_t^{-1}(x)$$

Note that $Slv (h, t) l^2 d\mu_t = Slv (T_t(y), t) l^2 d\mu_0(y)$ $= SIT(y) - y l^2 d\mu_0(y)$ $= W^2 (\mu_0, \mu_1)$ Note that, by defn of T_t and v $T_t := ((1-t)id + tT),$

 $d_{t}T_{t} = T - id$ $= v(T_{t}, t)$ Thus, $M_t := T_t \# \mu_0$, is a weak solution of (CE). \square

Now, we will use this PDE characterization of constant speed geodesics to characterize abs cts curves in (P2(IR^d), W2) and prove Benamon-Brenier thm.

We will do this via approximating abs cts curves by (p.w.) geodesics, O and we will need a way to conclude the limit is also a solu to $((\overline{E}))$. m=5je Def: If u e P(Rd), me Ms(Rd), upmee Id, the kinetic energy is $\mathbb{B}(\mu,m) = \frac{1}{2} \int \frac{|m(x)|^2}{\mu(x)} dx - S|v|^2 d\mu$ How to extend to general MERRA)? Def: For $r \in \mathbb{R}$, $\chi \in \mathbb{R}^{d}$ $\int_{\mathbb{B}}^{1} |\chi|^{2} \int_{1}^{1} |\chi|^{2} |\chi|^{2} \int_{1}^{1} |\chi|^{2} |\chi|^{2} |\chi|^{2} \int_{1}^{1} |\chi|^{2} |$ $if r^{>}0$ $if r = \chi = 0$ $ifr=0, x^{\neq 0}$ or r<0.

 $\begin{array}{l} \underbrace{\mathcal{D}ef(abscts):Given \ \alpha \ metric \ space \ (X,d)}_{\chi: \ (0,T) \rightarrow X \ is \ absolutely \ ds, \ denoted \ \chi \in \mathcal{AC}(0,T;X)_{1} \ if \ \exists \ g \in L^{2}(\{0,T\}) \ s.t. \end{array}$ $d(x(t_0), x(t_0)) \in \int_{t_0}^{t_0} g(s) ds \quad \forall \ 0 < t_0 \in t_0 \in T.$

 $R_{m,k}$: if this holds for $q(s) \equiv C \in \mathbb{R}$, then $\chi(t)$ is Lipschitz continuous.

ina vector space

Fact (Rademacher Thm): Given XEAC(O,T;X) tx'llt) exists for a.e. t and

$$d(x(t_0), x(t_0)) \leq \int_{t_0}^{t_0} |x'|(s)ds \quad \forall \quad 0 < t_0 \leq t_0 \leq t_0 \leq t_0$$

Theorem (characterization of abocts; curves in W2
(i) Suppose
$$\mu: [0,T] \rightarrow B_2(\mathbb{R}^d)$$
 is abjcts.
Then I to s.t. (μ, ν) is a dist som of cty
eqn and $\|\nu(\cdot, t)\|_{L^2(\mu t)} \leq |\mu|(t) \text{ a.e. } t \in [0,T].$
I kinetic energy time t

PL: "=>" Will prove in special case that ult) is concentrated on a cpt set KERQ. WL06 T=1 Strategy construct a sequence of simple solut to cty equilibriant converge to u. Let le be a mollifier with M2(q1)2+00

Given kell, consider discrete time sequence $\mu(^{\prime}k), \mu(^{\prime}k), ..., \mu(^{\prime}k), ..., \mu(^{k}k)$ and the mollified sequence $\mu_{^{\prime}k}^{^{k}} := P_{^{\prime}k} * \mu(^{^{\prime}}k), i=0,...,k$

B/c this sequence is << Id, the geodesics between each two elements of the sequence are sons to cty eqn. Chaining these geodesics toyether, we get a', dist som af cty eqn (µ, v) reparametrize reparametrige time su 1. [0,] -> P2(Rd)

 $\|\nabla^{k}(\cdot,t)\|_{L^{q}(\mathbb{N}_{+}^{k})}^{2} = K^{2} W_{2}^{2}(\mu_{i,k}^{k} \mid \mu_{i+1/k}^{k}) \text{ for } t \in [\frac{1}{k}, \frac{i+1}{k}]$ $\leq k^2 W_2^2(\mu(i/k),\mu(i/k))$ $\leq \left(k \int |\mu'|(s)ds\right)^{2} \int Jensen's$ $\leq k \int |\mu'|^{2}(s)ds$

Define m^k(t):= v^k(., t)du^k(t), which is a vector measure on TR² concentrated on a compact set.

Furthermore,

$$Im \#(t)(IRQ) = \left(\int_{RQ} |v + (-, t)| d\mu_{t}^{k} \right)^{2} \int_{RQ}^{1/2} \left(\int_{RQ} |v + (-, t)|^{2} d\mu_{t}^{k} \right)^{1/2} \\ \leq \left(\int_{RQ} |v + (-, t)|^{2} d\mu_{t}^{k} \right)^{1/2} \\ \leq k \int_{L} |\mu^{1}|(s) ds \quad \text{for } t \in [\frac{1}{k}, \frac{1+1}{k}) \\ \int_{L} \int_{L}$$

In pantialan, since
$$M_{\ell}^{\ell}(\mathbb{R}^{d}) \equiv 1 \in L^{1}(0, T]$$
,
 $M^{\ell}dt \rightarrow \mu_{\ell}dt$ narrowly wrt $G(\mathbb{R}^{d} \sim [0, T])$.
Thus, $\forall \ \mathcal{Q} \in C_{c}^{\infty}(\mathbb{R}^{d} \times [0, T])$
 $\int_{0}^{T} \int_{\mathbb{R}^{d}} (\partial_{t} \mathcal{Q}_{k_{1},t}) + v^{t}(x,t) \cdot \nabla \mathcal{Q}_{k_{1},t}) d\mu^{t}dt = 0$
 $\int_{0}^{T} \mathbb{R}^{d} (\partial_{t} \mathcal{Q}_{k_{1},t}) + v^{t}(x,t) \cdot \nabla \mathcal{Q}_{k_{1},t}) d\mu^{t}dt = 0.$
Since $M_{t}^{k} = v^{k}_{t} d\mu^{k}_{t}$, $m^{k} < \mu^{k}_{t}$, and
 $B(\mu^{k}_{t_{1}}, m^{k}_{t}) = \int |v^{k}|^{2} d\mu^{k}_{t}$
 $i^{t}(x)$